

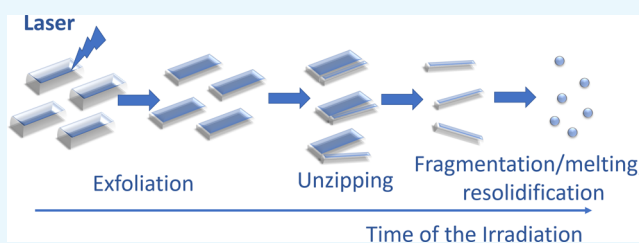
# Unveiling the Structure of MoS<sub>x</sub> Nanocrystals Produced upon Laser Fragmentation of MoS<sub>2</sub> Platelets

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## Supporting Information

**ABSTRACT:** Transition-metal dichalcogenide MoS<sub>2</sub> nanostructures have attracted tremendous attention due to their unique properties, which render them efficient nanoscale functional components for multiple applications ranging from sensors and biomedical probes to energy conversion and storage devices. However, despite the wide application range, the possibility to tune their size, shape, and composition is still a challenge. At the same time, the correlation of the structure with the optoelectronic properties is still unresolved. Here, we propose a new method to synthesize various morphologies of molybdenum sulfide nanocrystals, on the basis of ultrashort-pulsed laser fragmentation of MoS<sub>2</sub> platelets. Depending on the irradiation conditions, multiple MoS<sub>x</sub> morphologies in the form of nanoribbons, nanospheres, and photoluminescent quantum dots are obtained. Besides the detailed structural analysis of the various crystals formed, the structure–property relation is investigated and discussed.



## 1. INTRODUCTION

Transition-metal dichalcogenides (TMDs) are unique two-dimensional (2D) materials due to their fascinating electronic, optical, and catalytic properties. The MoS<sub>2</sub>, in particular, is a typical TMD comprising weakly bonded S–Mo–S sandwiched layers.<sup>1</sup> A single MoS<sub>2</sub> layer exhibits remarkably different properties than those of the bulk; the most pronounced one is that it is a direct-band gap semiconductor<sup>2,3</sup> contrary to the bulk, which shows an indirect band gap of 1.22 eV.<sup>4</sup> Their unique properties render them efficient nanoscale functional components for various applications, including sensors,<sup>5</sup> biomedicine,<sup>6</sup> and energy conversion and storage devices.<sup>7,8</sup>

Besides this, zero-dimensional (0D) TMDs have been found to exhibit high stability and extraordinary optoelectronic properties due to quantum confinement effects.<sup>9–11</sup> Such low-dimensional materials found numerous applications in electrochemical photocatalysis,<sup>12,13</sup> bioimaging,<sup>14,15</sup> and optoelectronic devices.<sup>16</sup> MoS<sub>2</sub> nanostructures, in particular, have been synthesized with different methods, such as ionic liquid-assisted grinding exfoliation,<sup>12</sup> hydrothermal routes,<sup>10</sup> electrochemical exfoliation,<sup>17</sup> colloidal methods at high temperature,<sup>18</sup> laser ablation in liquids,<sup>19–26</sup> and via an induced Fenton reaction.<sup>27</sup> Among them, the grinding exfoliation<sup>12</sup> and the hydrothermal<sup>10</sup> methods are time-consuming, considering that the first requires multiple centrifugation steps for the size selection of the particles whereas the second demands 18 h of reaction at 220 °C. Besides this, the electrochemical methods are two-step procedures and use hazardous materials.<sup>17,27</sup> Finally, the wet chemistry synthesis requires complex apparatus such as a Schlenk line and continuous inert gas flow during the

reaction and takes place at high temperature, which increases the time and cost of the final products.<sup>18</sup>

The shape of the MoS<sub>2</sub> nanocrystals synthesized via the methods developed to date, is centrosymmetric, whereas the possibility to tune their size, shape, and composition is still a challenge. In this work, we propose a simple and rapid photoinduced method to control the morphology and stoichiometry of MoS<sub>x</sub> nanostructures, via ultrashort-pulsed laser irradiation of MoS<sub>2</sub> platelets, in solution. It is a single-step, room temperature, and green process, providing different nanocrystal morphologies without the need of hazardous materials and additional surface functionalization. Besides this, it allows fundamental studies of the impact of crystal structure and morphology on the physicochemical properties of the MoS<sub>x</sub> nanostructures. The rather simple and scalable method under ambient conditions provides unique opportunities for the cost-effective synthesis of bulk amounts of nanostructures with controllable size, shape, and stoichiometry.

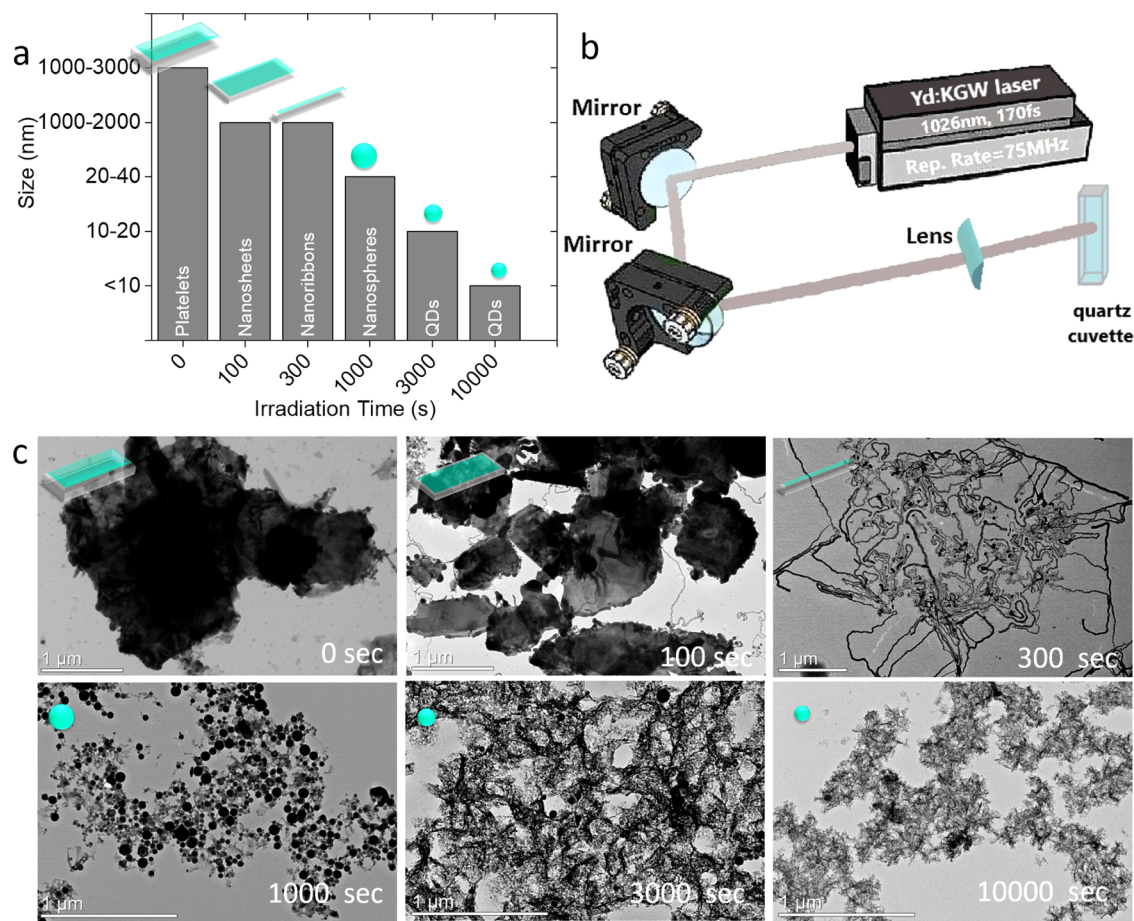
## 2. RESULTS AND DISCUSSION

The initial colloidal solution of platelets has been prepared by laser ablation of a MoS<sub>2</sub> powder pellet in deionized water. For this purpose, a 513 nm laser source, emitting 170 fs pulses with a repetition rate of 60 kHz, has been used (Figure S1). The platelet colloids obtained by this process were stable in solution, and no precipitation was observed upon storage for months at ambient conditions. Subsequently, the aqueous

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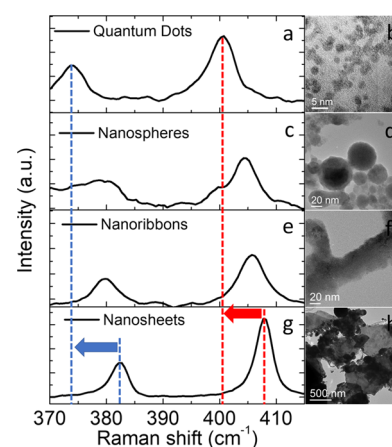


**Figure 1.** Size and shape evolution of the MoS<sub>x</sub> structures by increasing the irradiation time of the platelet aqueous liquid (a). Schematic representation of the experimental setup utilized for the photoinduced fabrication of the nanostructures in liquid media under ambient conditions (b). Low-magnification transmission electron microscopy (TEM) images of the structures (c).

solution of MoS<sub>2</sub> platelets was irradiated using a 1030 nm laser source, emitting 170 fs pulses with a high repetition rate of 75 MHz (Figure 1b).

The influence of the irradiation time, ranging from 0 to 10 000 s, on the morphological and structural evolution of the MoS<sub>x</sub> nanostructures attained (Figure 1a) was carefully examined by TEM and Raman experiments. Representative, bright-field low-magnification TEM images for all structures obtained are presented in Figure 1c. In particular, the initial colloidal solution before the irradiation comprises bulk-like platelets, which have been ablated from the MoS<sub>2</sub> target. Following 100 s of irradiation, the as-prepared platelets become thinner (nanosheets), as indicated by the lighter contrast exhibited in the respective TEM images, while their lateral size becomes smaller. At 300 s of irradiation, nanoribbons of 20–50 nm width and length similar to that of the nanosheets prepared at 100 s are formed. Finally, above 1000 s of irradiation time, nanospheres are formed, with sizes down to 5 nm; the higher the exposure time, the lower is the mean nanosphere diameter produced.

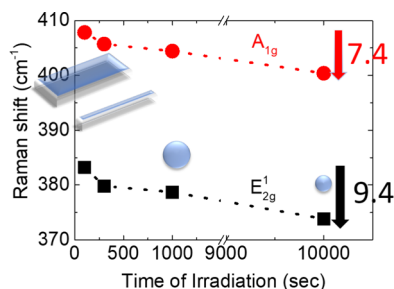
Raman spectroscopy was used to probe the vibrational properties of the different structures formed. As shown in Figure 2, the two Raman peaks associated with <sup>1</sup>E<sub>2g</sub> (in-plane motion of Mo and S in opposite directions) and A<sub>1g</sub> (out-of-plane motions of S atoms) active modes are observed in all cases. In particular, for the nanosheets formed at 100 s, the two peaks are centered at 383.2 and 407.8 cm<sup>-1</sup> (Table S1), in



**Figure 2.** Raman spectra and TEM images: quantum dots obtained at 10 000 s of irradiation (a, b), nanospheres at 1000 s (c, d), nanoribbons at 300 s (e, f), and nanosheets at 100 s (g, h).

accordance to previous reports on bulk MoS<sub>2</sub>.<sup>28</sup> Upon transition from nanosheets to nanoribbons and nanospheres, a characteristic blue shift and broadening of both peaks were observed. The blue shift is 3.4 cm<sup>-1</sup> for the <sup>1</sup>E<sub>2g</sub> and 2.1 cm<sup>-1</sup> for the A<sub>1g</sub> for the nanosheet-to-nanoribbon transition. According to the literature, the thinning of bulk MoS<sub>2</sub> to ultrathin nanosheets comprising three layers leads to a blue shift of the <sup>1</sup>E<sub>2g</sub> and A<sub>1g</sub> modes of about 0.6 and 2.1 cm<sup>-1</sup>,

respectively.<sup>28</sup> In addition, the broadening of these peaks has been associated with the decrease in the number of layers.<sup>29</sup> On the basis of these findings, the large blue shift of the  $^1E_{2g}$  mode could not be solely attributed to the nanostructures thinning; therefore, another structural parameter could also play a role (Figure 3). No significant shift was observed, for



**Figure 3.** Peak position of the  $^1E_{2g}$  and  $A_{1g}$  Raman modes, extracted by fitting the peaks in the Raman spectra with a Gaussian curve, for the different morphologies obtained ranging from platelets to quantum dots.

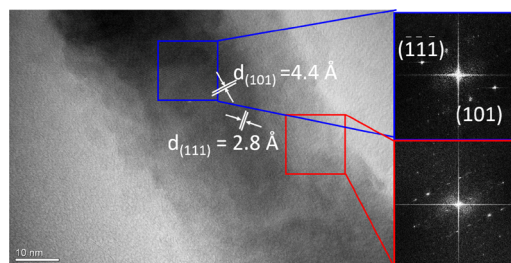
both modes, upon the transition from nanoribbons to nanospheres. However, the transformation of nanospheres to quantum dots gives rise to a further blue shift for both modes, namely, 4.9 for the  $^1E_{2g}$  and 4  $\text{cm}^{-1}$  for the  $A_{1g}$ , respectively. This observation is in contrast to the majority of literature results on Raman spectra of  $\text{MoS}_x$  nanostructures. For example, a red shift of the in-plane mode has been observed in monolayered structures of nanosheet<sup>28</sup> or quantum dot<sup>30</sup> morphologies. A large blue shift has been found only in the case of quantum dots prepared with an ionic liquid-assisted grinding exfoliation process, associated with disorder at the particles edges.<sup>12</sup> In addition, such large shift of the  $^1E_{2g}$  mode ( $2.1 \text{ cm}^{-1}$  per % strain) has been found in strained  $\text{MoS}_2$  structures.<sup>31</sup> The shift is larger in the case of few-layered structures ( $2.1 \text{ cm}^{-1}$  per % strain) compared with monolayered ones ( $1.7 \text{ cm}^{-1}$  per % strain).

Furthermore, there are three different mechanisms that would explain the asymmetric broadening of the Raman peaks: (1) The broadening in the nanocrystalline materials is described by the phonon confinement model,<sup>32</sup> which leads to a breakup of selection rules, meaning, also, non-zone-center phonons will participate in scattering.<sup>33–35</sup> (2) The size/shape of the spectral lines is changing due to size effects of the grown nanostructures.<sup>36</sup> (3) The crystallinity of the material could affect the broadening of the peaks. A well-crystalline material yields a spectrum with very sharp, intense Raman peaks, whereas an amorphous material will show broader less intense Raman peaks. These two states (e.g., fully amorphous or fully crystalline) can be considered as spectral extremes, and a Raman spectrum from an intermediate state (e.g., partially crystalline) will have characteristics that are intermediate in terms of peak intensity and width (sharpness). On the basis of the Raman spectroscopy findings on disordered particles and strained layered structures and the three main mechanisms that could affect the broadening of the peaks stated above, we may postulate that apart from the reduction in the number of atomic layers, the laser-induced structural changes give rise to remarkable strain and disorder effects.

High-resolution transmission electron microscopy (HRTEM) analysis was used to evaluate the formation

mechanism of the various kinds of nanostructures obtained. In particular, the following was observed.

**2.1. For nanoribbons.** HRTEM imaging revealed that the nanoribbons prepared upon 300 s of irradiation are crystalline (Figure 4) and have a width of  $w = 31.6 \pm 11.1 \text{ nm}$  (Figure



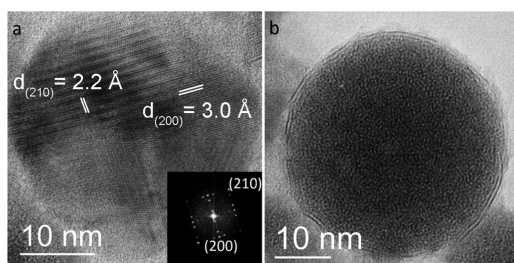
**Figure 4.** HRTEM image of an individual nanoribbon. The corresponding FFT patterns for the selected areas marked by rectangles in the middle (blue) and at the edge (red) of the nanoribbon are also presented. The zone axis in the image corresponds to the  $[101]$  direction. The indexing has been assigned according to the reference pattern of the monoclinic  $\text{Mo}_2\text{S}_3$  structure (ICSD, #073453).

S2a). The corresponding FFT patterns for the selected area, marked by a blue rectangle in the middle of the nanoribbons in Figure 4, cannot be indexed with the hexagonal structure of the  $\text{MoS}_2$  (ICSD, #049801). The structure is compatible with the monoclinic structure of the  $\text{Mo}_2\text{S}_3$  (ICSD, #073453), as the measured lattice spacings for the (111) and (101) are 2.8 and 4.4 Å, respectively. The corresponding  $\text{Mo}_2\text{S}_3$  polymorph has been found in the literature to be formed due to surface phase transformation upon thermal annealing of  $\text{MoS}_2$  at a temperature up to 1300 K<sup>37</sup> or by sputtering of  $\text{MoS}_2$  films with low-energy  $\text{Ar}^+$  ions<sup>38</sup>. Besides this, irradiation of  $\text{MoS}_2$  monolayers with an electron beam leads to the introduction of sulfur vacancies.<sup>39</sup> Finally, the  $\text{MoS}_x$  phase was found in nanodots synthesized by the arc-melting method at 3000 °C, followed by a ball milling process.<sup>40</sup>

To further shed light on the structural features of the nanoribbons, we compared the FFT pattern calculated from an area in the middle to that at the edge of an individual nanoribbon (Figure 4). Although the central part is well-crystalline and there is no evidence for grain boundaries, the FFT pattern corresponding to the nanoribbon edge revealed the presence of disorder; the observed streaking in diffraction spots, in particular, indicates the existence of a stacking fault area along the  $b$  axis.

**2.2. For nanospheres.** Following 1000 s of laser irradiation, nanospheres of  $d = 36.2 \pm 10.9 \text{ nm}$  in diameter (Figure S2b) have been formed. HRTEM analysis revealed the presence of different types of particles. The majority are single-phase twinned particles (Figure 5a), and a relatively small number of fullerene-like or amorphous particles (Figures 5b and 5c) has been additionally observed. Fullerene particles have been synthesized by a similar laser-ablation technique in aqueous solution.<sup>26,41</sup> FFT pattern analysis revealed the crystal structure of the single-phase particles, which is compatible with the monoclinic structure of the  $\text{Mo}_2\text{S}_3$  (ICSD, #073453). Indeed, lattice fringes with interplanar spacings of 3.0 and 2.2 Å have been identified, which are assigned to the (200) and (210) of the monoclinic crystal  $\text{Mo}_2\text{S}_3$  structure, respectively. It should be noted here that the crystal structure of such nanocrystals is the same as that of the nanoribbons, although a



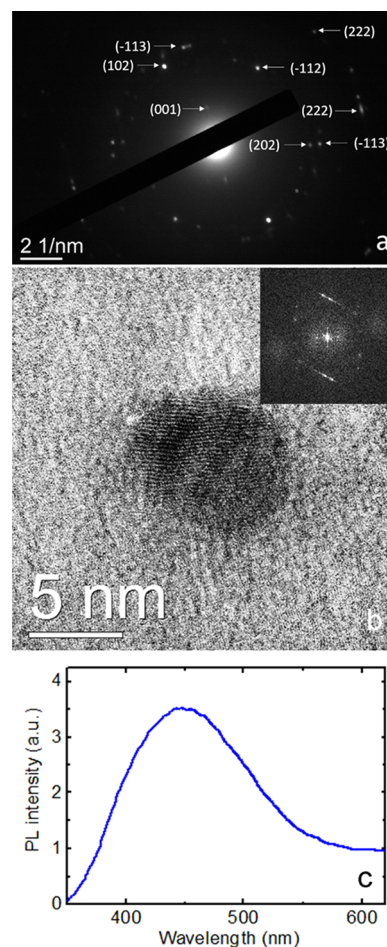


**Figure 5.** HRTEM images from the two different types of nanospheres, namely, single-phase twined (a) and fullerene-like (b) nanosphere. The inset in (a) is an FFT pattern corresponding to an individual nanosphere. The zone axis is the [001]. The indexing has been assigned according to the reference pattern of the monoclinic structure of the  $\text{Mo}_2\text{S}_3$  (ICSD, #073453).

small percentage of nanoparticles exhibit the rhombohedral crystal symmetry of the  $\text{Mo}_3\text{S}_4$  (Figure S3). The coexistence of these two phases in the same sample has been also found in previous reports.<sup>40</sup> The mixture of several molybdenum multisulfides has also been confirmed from the stoichiometry found by energy-dispersive spectrometry (EDS) analysis in different areas of the sample (Figure S4). This stoichiometry was found to be MoS (1:1), which is far away from the stoichiometry Mo/S = 1:2 of the  $\text{MoS}_2$  target. Upon careful inspection of the HRTEM images and the corresponding FFT pattern, it can be concluded that the particles exhibit a quasispherical morphology of the twined type, revealed by the additional spots between the spots of the monoclinic phase. Variation of the Mo/S ratio in  $\text{MoS}_2$  crystal structures has been found in the literature that it is associated with the presence of S vacancies.<sup>42</sup> Twin defects, dislocations, and grain boundaries in the crystal structure are accumulated by sulfur vacancies.<sup>43,44</sup> A ratio of 1.2 found in substoichiometric bilayers of  $\text{MoS}_2$  due to lattice distortion originated from exfoliation process using hydrazine salts.<sup>45</sup>

**2.3. For the quantum dots.** Quantum dots with diameter  $d = 6.5 \pm 1.6$  nm (Figure S2c) are formed above 10 000 s of laser irradiation. Selected area diffraction analysis of a random area of the quantum dot sample (Figure S6) revealed the crystal structure of the quantum dots, which is compatible to the monoclinic crystal structure of  $\text{Mo}_2\text{S}_3$  (Figure 6a). No indication for molybdenum oxide has been found due to possible oxidation of the surface.<sup>46–48</sup> This conclusion is in good agreement with the Raman spectra at the higher shift region. Peaks at 666 and 820 nm that could be assigned to the stretching modes of the triply coordinated oxygen ( $\text{Mo}_3\text{--O}$ ) and the doubly coordinated oxygen ( $\text{Mo}_2\text{--O}$ ) of the  $\alpha\text{-MoO}_3$  phase<sup>49,50</sup> are found only for the case of the nanospheres (Figure S7). The corresponding FFT pattern calculated from the HRTEM images (Figures 6b and S8) indicated a twined structure.

Notably, the quantum dots are photoluminescent at room temperature upon their excitation by a 248 nm laser beam (Figure 6c). In particular, a single PL peak is observed centered at 448 nm whereas no additional peaks are observed at longer wavelengths, indicating the absence of quantum dot aggregation. A similar spectrum has been recorded for the nanospheres (Figure S9). The main features are similar, but the spectrum of the nanospheres is blue-shifted compared to that of quantum dots. This is a contradictory result if we consider the particle size of the two samples. This is due to

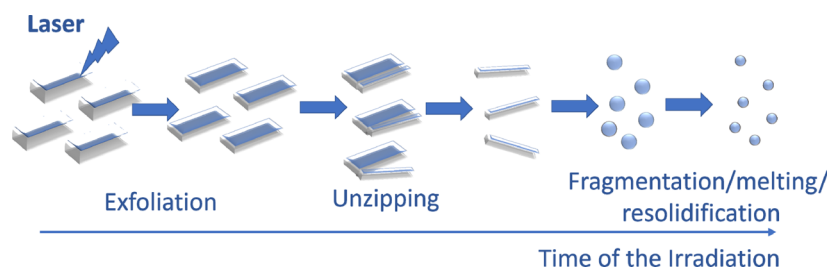


**Figure 6.** Selected area diffraction pattern (SAED) (a), HRTEM image (b), and photoluminescence spectrum (c) from the quantum dots formed above 10 000 s of laser irradiation. The inset in (b) is the corresponding FFT pattern calculated for an individual quantum dot (right). The indexing of the SAED has been assigned according to the reference pattern of the monoclinic structure of the  $\text{Mo}_2\text{S}_3$  (ICSD, #073453).

existence of different particles (mainly of  $\text{MoO}_3$ ) in the sample of the nanospheres.

Similar spectra have been recorded for water soluble 1,4-diaminobutane-functionalized  $\text{MoS}_2$  quantum dots of about 5 nm synthesized by a bottom-up approach.<sup>51</sup>

On the basis of the above analysis, we may speculate on the formation mechanism that accounts for the structure and shape evolution of the various  $\text{MoS}_x$  nanostructures observed (Figure 7). Initially, laser irradiation of  $\text{MoS}_2$  platelets gives rise to the photoinduced exfoliation to form few-layered nanosheets. The subsequent irradiation of nanosheets leads to the photo-induced unzipping<sup>52–54</sup> of the nanosheets' edges to form nanoribbons, which are crystalline in the  $\text{Mo}_2\text{S}_3$  structure, exhibiting structural defects at their edges. This explanation complies with previous reports, where the  $\text{Mo}_2\text{S}_3$  polymorph was found at the edges of  $\text{MoS}_2$  nanosheets.<sup>37</sup> Upon further irradiation, the nanoribbons transform in nanospheres through a laser-induced melting and subsequent resolidification.<sup>41,55</sup> The temperature in the irradiated area due to the laser energy absorption is increased above the melting point of the material, and melted  $\text{MoS}_x$  droplets are formed. The resolidification occurs when the surface temperature of the droplets reduces below the melting point of  $\text{MoS}_x$ . The preferable shape of



**Figure 7.** Photoinduced formation mechanism for the different MoS<sub>x</sub> morphologies.

these structures tends to be spherical to reduce their surface energy. The observed strained structure of the nanoparticles revealed by HRTEM analysis could explain the large blue shift observed in the respective Raman modes. Finally, the nanoparticles break into smaller particles in the form of quantum dots, due to laser-induced ablative fragmentation of the larger nanospheres<sup>56,57</sup> indicated by the preservation of the initial Mo<sub>2</sub>S<sub>3</sub> twined structure.

### 3. CONCLUSIONS

In summary, this work illustrates that laser irradiation of MoS<sub>2</sub> bulklike platelets could be used as an effective way for the synthesis of different morphologies of molybdenum multi-sulfides. This approach is simple, scalable, performed under ambient conditions, and can therefore be useful for the cost-effective synthesis of bulk amounts of nanostructures with controllable size, shape, and stoichiometry. More important, the various nanocrystal morphologies can be produced without the need of hazardous chemicals and additional surface functionalization processes, which make this technique appealing for many potential applications.

### 4. MATERIAL AND METHODS

**4.1. Preparation of the Different Colloids.** **4.1.1. Production of the Aqueous Solution of MoS<sub>2</sub> Platelets.** Initially, a MoS<sub>2</sub> target has been prepared via pressing MoS<sub>2</sub> powder into a pellet of 15 mm diameter. This target was placed at the bottom of a glass vial and immersed into 2 mL of deionized water. Subsequently, the beam of a Yb:KGW femtosecond laser was focused through the water layer, onto the target to exfoliate the MoS<sub>2</sub> powder into platelets (Figure S1). The laser wavelength, the repetition rate, the pulse duration, and the power used for this purpose were 513 nm, 60 kHz, 170 fs, and 90 mW, respectively. The total ablation time was 10 min. During laser irradiation, the target was continuously rotated using a rotation stage. All solutions obtained after irradiation were centrifuged at 1500 RPM for 10 min. Then, the supernatant colloidal solution is placed in a quartz cuvette for subsequent irradiation.

**4.1.2. Synthesis of MoS<sub>x</sub> Nanostructures (Figure 1).** The as-prepared aqueous solution of the MoS<sub>2</sub> platelets was divided into 5 equivalent aliquots. Then, each aliquot was irradiated for 100, 300, 1000, 3000, and 10 000 s, respectively, using the same laser system, although at different irradiation parameters. In this case, the wavelength, repetition rate, pulse duration, and power used were 1026 nm, 75 MHz, 170 fs, and 4 mW, respectively. During the irradiation process, the colloidal solution was continuously agitated via use of a magnetic stirrer. The synthesis is reproducible, but the yield is low.

**4.2. Characterization of the Colloidal Solutions.**

**4.2.1. Transmission Electron Microscopy (TEM).** Low-

magnification and high-resolution TEM (HRTEM) images were recorded on a JEOL 2100 transmission electron microscope, operating at an accelerating voltage of 200 kV. For the purposes of the TEM analysis, a drop of the aqueous colloidal solution was deposited onto a carbon-coated copper TEM grid and then the solvent was allowed to evaporate. All images were recorded by a Gatan ORIUS SC 1000 CCD camera, whereas the structural features of the nanostructures were studied by two-dimensional (2D) fast Fourier transform (FFT) analysis.

**4.2.2. Raman Spectroscopy.** Raman measurements were performed at room temperature using a Horiba LabRAM HR Evolution confocal micro-spectrometer, in backscattering geometry (180°), equipped with an air-cooled solid-state laser operating at 532 nm, with 100 mW output power. The laser beam was focused on the samples using a 50× Olympus microscope objective (numerical aperture of 0.50), providing a ~22.5 mW power on each sample. Raman spectra over the 10–3000 cm<sup>-1</sup> wavenumber range (with an exposure time of 20 s and 5 accumulations) were collected by a Peltier cooled CCD (1024 × 256 pixels) detector at -60 °C, with a resolution better than 1 cm<sup>-1</sup> achieved thanks to an 1800 grooves/mm grating and an 800 mm focal length. Test measurements were carried out using different optical configurations, exposure time, beam power, and accumulations to obtain sufficiently informative spectra using a confocal hole of 100 μm but ensuring to avoid alteration of the sample, whereas the high spatial resolution allowed us to carefully verify the sample homogeneity. The wavelength scale was calibrated using a Silicon standard (520.7 cm<sup>-1</sup>), and the acquired spectra were compared with scientific published data and reference databases, such as Horiba LabSpec 6 (Horiba).

**4.2.3. Laser-Induced Fluorescence (LIF) Spectroscopy.** Each solution was placed in quartz cuvette without further dilution, and the respective PL spectra were recorded at room temperature. For sample excitation, a KrF excimer laser, operating at 248 nm has been utilized. The pulse duration was about 20 ns, the excitation energy ~3.2 mJ, and the laser beam diameter 9.0 mm (fluence ~5 mJ/cm<sup>2</sup>). The fluorescence measurements were performed at room temperature and recorded by a monochromator (PTI Technology) that is connected with an Andor iStar 734 Series Intensified Charge Coupled Device. The fluorescence signal of the samples was collected and guided to the spectrograph by an optical fiber.

**4.2.4. Scanning Electron Microscopy (SEM) and Energy-Dispersive X-ray Spectroscopy (EDS).** SEM (JEOL 7000), equipped with an EDS (INCA PentaFET-x3), was used for the analysis of the samples. These experiments were performed on solid-state samples produced after drying a small volume of the corresponding colloidal solution on a piece of a single-crystalline silicon wafer.

## ■ ASSOCIATED CONTENT

### ■ Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsomega.8b01390.

EDS analysis, TEM images, size distributions, Raman spectra, PL curves and scheme of the setup for the laser ablation of the flakes in liquid media (PDF)

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### Author Contributions

The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

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### Notes

The authors declare no competing financial interest.

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